

# Toward a self-consistent and unitary nuclear reaction network

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# Outline

- Overview: LANL light nuclear reaction program
- Unitarity: SBBN & beyond, reaction networks, R-matrix,  $^{17}\text{O}, ^9\text{B}$  examples
- Recent related development: EFT  $\longleftrightarrow$  R-matrix ( $a_c \rightarrow 0 \quad \forall c$ )
- Future work & conclusion

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# Light nuclear reaction program @ LANL

## ■ Motivation

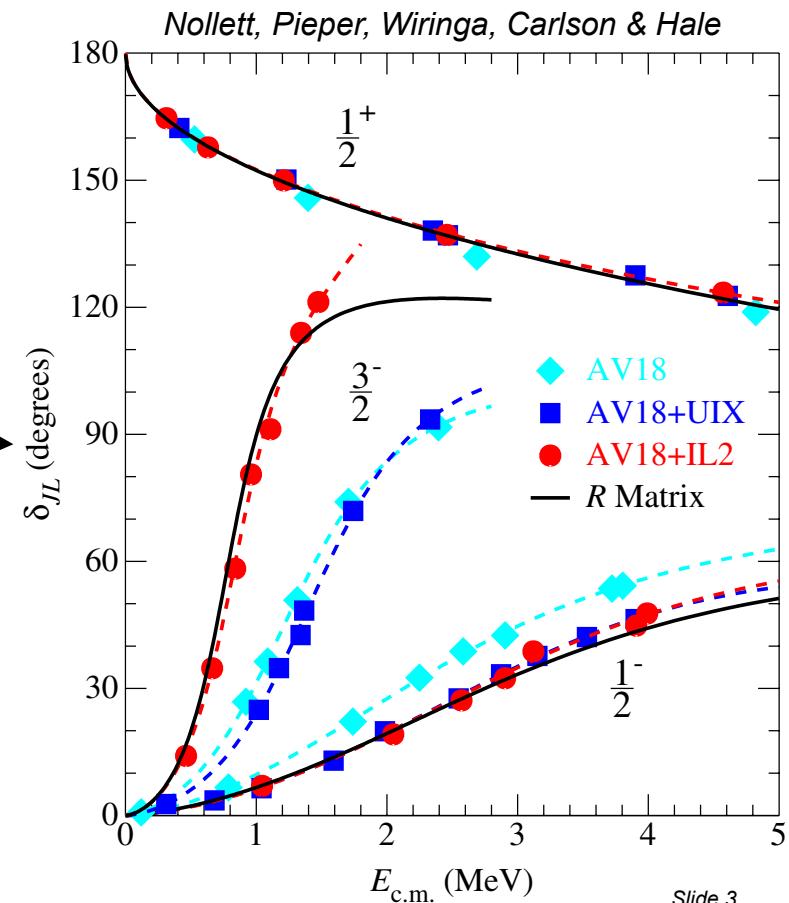
- Data sets:  $\sigma$ ,  $\sigma(\theta)$ ,  $A_i(\theta)$ ,  $C_{i,j}$ ,  $K_i^j$ ,  $\Sigma(\gamma)$ , ... → T matrix → resonance spectrum
- **Unitary** parametrization of compound nuclear system
- Applications: astrophysical, nuclear security, criticality safety, charge-particle transport, nuclear data (ENDF, ENSDF)

## ■ Ab initio

- Variational MC; Green's function MC
- GFMC [PRL 99, 022502 (2007)]
  - n- $^4$ He phase shifts
  - comparison GFMC/R-matrix
- challenge: multichannel (eg. na → dt)

## ■ Phenomenology

- R matrix (2 → 2 body scatt/reacs)
- 3-body: isobaric models, sequential decay



# EDA Analyses of Light Systems

A	System	Channels	Energy Range (MeV)
2	N-N	p+p; n+p, $\gamma$ +d	0-30 0-40
3	N-d	p+d; n+d	0-4
	$^4\text{H}$ $^4\text{Li}$	n+t p+ $^3\text{He}$	0-20
4	$^4\text{He}$	p+t n+ $^3\text{He}$ d+d	0-11 0-10 0-10
	$^5\text{He}$	n+ $\alpha$ d+t $^5\text{He}+\gamma$	0-28 0-10
5	$^5\text{Li}$	p+ $\alpha$ d+ $^3\text{He}$	0-24 0-1.4

# Analyses of Light Systems, Cont.

A	System (Channels)
6	$^6\text{He}$ ( $^5\text{He} + \text{n}$ , $\text{t} + \text{t}$ ); $^6\text{Li}$ ( $\text{d} + ^4\text{He}$ , $\text{t} + ^3\text{He}$ ); $^6\text{Be}$ ( $^5\text{Li} + \text{p}$ , $^3\text{He} + ^3\text{He}$ )
7	$^7\text{Li}$ ( $\text{t} + ^4\text{He}$ , $\text{n} + ^6\text{Li}$ ); $^7\text{Be}$ ( $\gamma + ^7\text{Be}$ , $^3\text{He} + ^4\text{He}$ , $\text{p} + ^6\text{Li}$ )
8	$^8\text{Be}$ ( $^4\text{He} + ^4\text{He}$ , $\text{p} + ^7\text{Li}$ , $\text{n} + ^7\text{Be}$ , $\text{p} + ^7\text{Li}^*$ , $\text{n} + ^7\text{Be}^*$ , $\text{d} + ^6\text{Li}$ )
9	$^9\text{Be}$ ( $^8\text{Be} + \text{n}$ , $\text{d} + ^7\text{Li}$ , $\text{t} + ^6\text{Li}$ ); $^9\text{B}$ ( $\gamma + ^9\text{B}$ , $^8\text{Be} + \text{p}$ , $\text{d} + ^7\text{Be}$ , $^3\text{He} + ^6\text{Li}$ )
10	$^{10}\text{Be}$ ( $\text{n} + ^9\text{Be}$ , $^6\text{He} + \alpha$ , $^8\text{Be} + \text{nn}$ , $\text{t} + ^7\text{Li}$ ); $^{10}\text{B}$ ( $\alpha + ^6\text{Li}$ , $\text{p} + ^9\text{Be}$ , $^3\text{He} + ^7\text{Li}$ )
11	$^{11}\text{B}$ ( $\alpha + ^7\text{Li}$ , $\alpha + ^7\text{Li}^*$ , $^8\text{Be} + \text{t}$ , $\text{n} + ^{10}\text{B}$ ); $^{11}\text{C}$ ( $\alpha + ^7\text{Be}$ , $\text{p} + ^{10}\text{B}$ )
12	$^{12}\text{C}$ ( $^8\text{Be} + \alpha$ , $\text{p} + ^{11}\text{B}$ )
13	$^{13}\text{C}$ ( $\text{n} + ^{12}\text{C}$ , $\text{n} + ^{12}\text{C}^*$ )
14	$^{14}\text{C}$ ( $\text{n} + ^{13}\text{C}$ )
15	$^{15}\text{N}$ ( $\text{p} + ^{14}\text{C}$ , $\text{n} + ^{14}\text{N}$ , $\alpha + ^{11}\text{B}$ )
16	$^{16}\text{O}$ ( $\gamma + ^{16}\text{O}$ , $\alpha + ^{12}\text{C}$ )
17	$^{17}\text{O}$ ( $\text{n} + ^{16}\text{O}$ , $\alpha + ^{13}\text{C}$ )
18	$^{18}\text{Ne}$ ( $\text{p} + ^{17}\text{F}$ , $\text{p} + ^{17}\text{F}^*$ , $\alpha + ^{14}\text{O}$ )

# LANL R-matrix light element standards: status

**n+p:** no new work since 2008; cross sections in pretty good shape below 30 MeV; main need is extension to higher energies (150-200 MeV), with associated covariances.

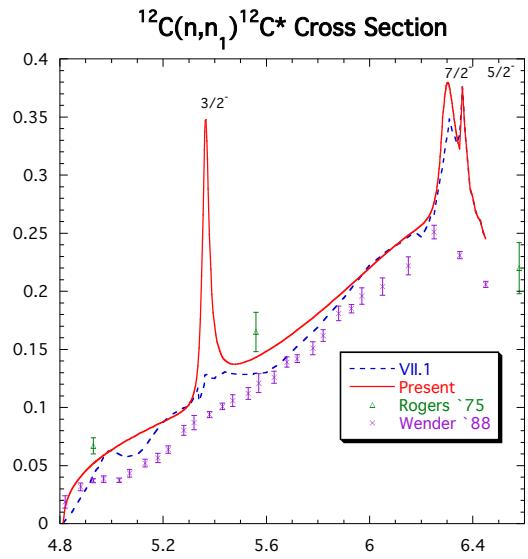
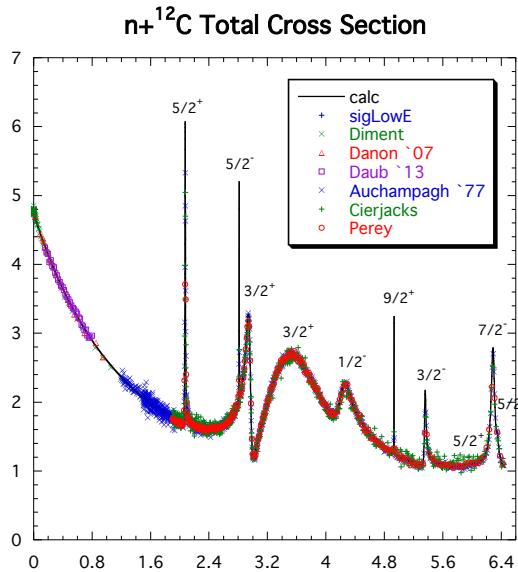
**n+<sup>3</sup>He:** Some new work, especially for n+<sup>3</sup>He capture; <sup>3</sup>He(n,n)<sup>3</sup>He scattering data re-worked by Drosd and Lisowski – could be used in a new analysis of the <sup>4</sup>He system.

**n+<sup>6</sup>Li:** Some new work on <sup>7</sup>Li system around 2008 included new LANSCE measurements of <sup>6</sup>Li(n,t)<sup>4</sup>He differential cross section – was included in ENDF/B-VII.1 above 1 MeV. Cross sections should be re-visited below 1 MeV, although there may not be any new data.

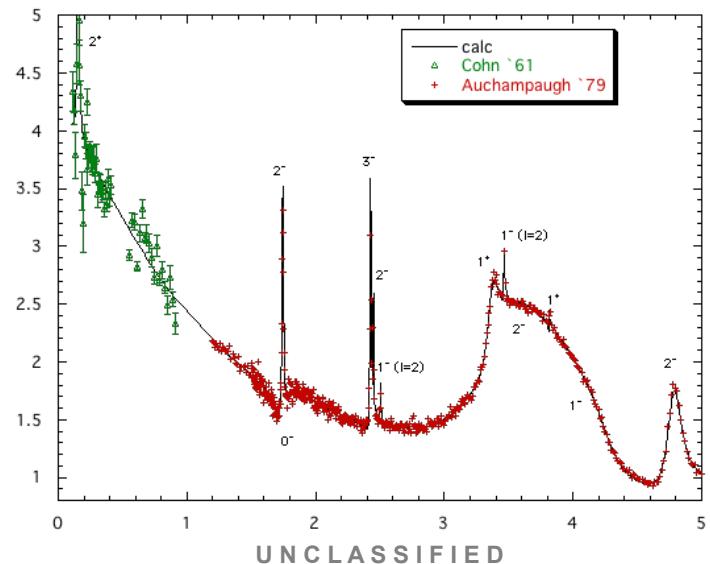
**n+<sup>10</sup>B:** No new work since last standards evaluation. New data from Geel, Ohio U. [including first measurement of the (n,p) cross section]. R-matrix analysis for <sup>11</sup>B system should be extended above 1 MeV.

**n+<sup>12,13</sup>C:** Considerable new work in the last couple of years. New data for <sup>12</sup>C(n,n'γ) from Geel, Los Alamos changed the (n,n') cross sections. Isotopic evaluations combined to make more accurate standard evaluation for C-0.

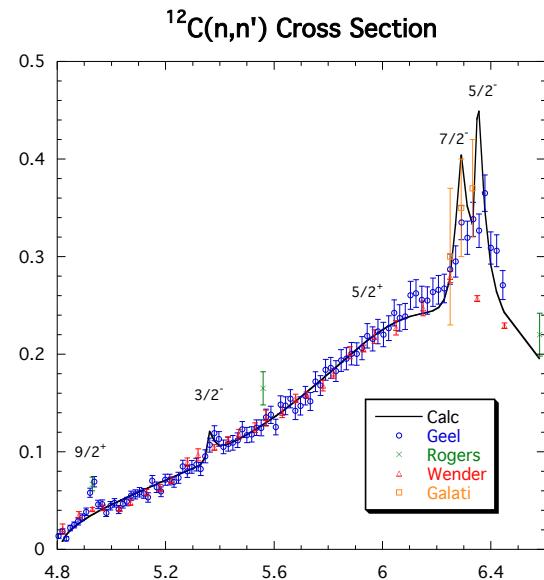
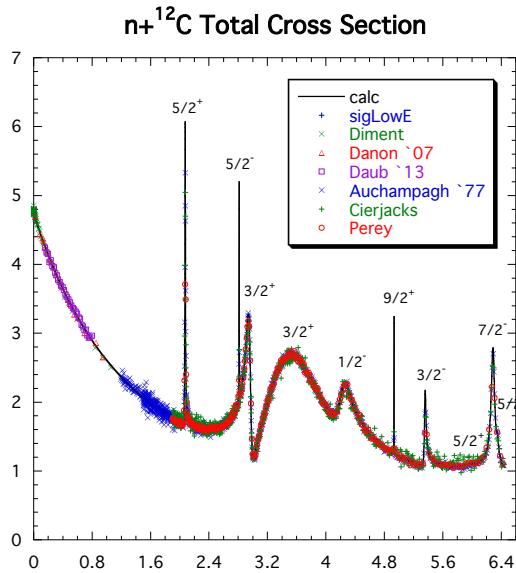
# $^{13,14}\text{C}$ system analyses: $\sigma_{\text{T}}$ (b) vs. $E_n$ (MeV)



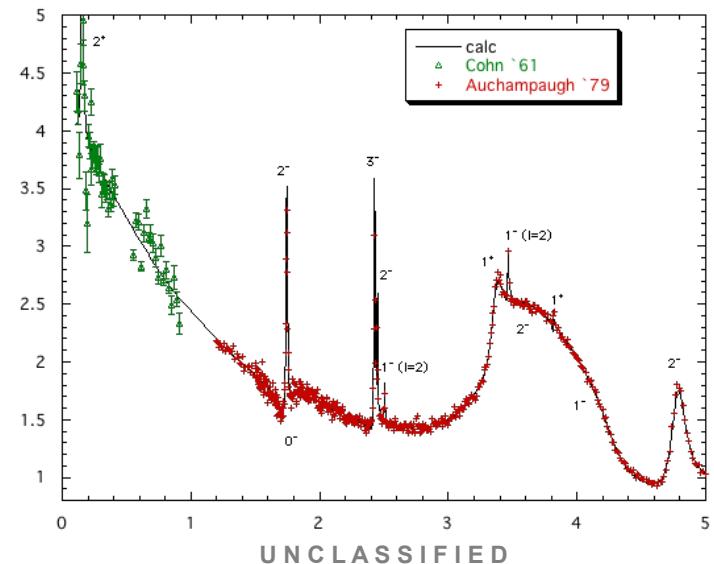
$n + ^{13}\text{C}$  Total Cross Section



# $^{13,14}\text{C}$ system analyses: $\sigma_{\text{T}}$ (b) vs. $E_n$ (MeV)



$n + ^{13}\text{C}$  Total Cross Section



# Unitary, self-consistent primordial nucleosynthesis

## ■ State of standard big-bang nucleosynthesis (BBN)

- d &  $^4\text{He}$  abundances: signature success cosmology+nucl astro+astroparticle
  - but there's at least one **Lithium ( $^7\text{Li}$ ) Problem** [ $^6\text{Li}$  too? *Lind et.al. 2013*]
- coming *precision* observations of d,  $^4\text{He}$ ,  $\eta$ ,  $N_{\text{eff}}$  demand new BBN capabilities
- resolution of  $^7\text{Li}$  problem:
  - observational/stellar astrophysics?
  - $^7\text{Li}$  controversial anomaly: nuclear physics solution?
  - new physics?

## ■ Advance BBN as a tool for precision cosmology

- incorporate **unitarity** into strong & electroweak interactions
- couple **unitary reaction network (URN)** to full Boltzmann transport code
  - neutrino energy distribution function evolution/transport code
  - fully coupled to nuclear reaction network
  - calculate light primordial element abundance for non-standard BBN
    - active-sterile  $\nu$  mixing
    - massive particle out-of-equilibrium decays→energetic active SM particles
- Produce tools/codes for nuc-astro-particle community: test new physics w/BBN
  - existing codes are based on Wagoner's (1969) code

# Formal unitarity: consequences

$$\left. \begin{array}{l} \delta_{fi} = \sum_n S_{fn}^\dagger S_{ni} \\ S_{fi} = \delta_{fi} + 2i\rho_f T_{fi} \\ \rho_n = \delta(H_0 - E_n) \end{array} \right\} \quad T_{fi} - T_{fi}^\dagger = 2i \sum_n T_{fn}^\dagger \rho_n T_{ni}$$

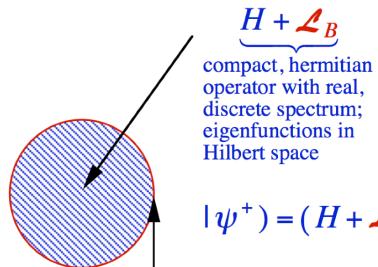
NB: **unitarity** implies optical theorem  $\sigma_{\text{tot}} = \frac{4\pi}{k} \text{Im } f(0)$ ; but *not only* the O.T.

## ■ Implications of **unitarity** constraint on transition matrix

1. Doesn't uniquely determine  $T_{ij}$ ; highly restrictive, however
  - Elastic:  $\text{Im } T_{11} = -\rho_1$  (assuming T & P invariance)
  - Multichannel:  $\text{Im } \mathbf{T} = -\boldsymbol{\rho}$
2. Unitarity violating transformations
  - cannot scale **any** set:  $T_{ij} \rightarrow \alpha_{ij} T_{ij} \quad \alpha_{ij} \in \mathbb{R}$
  - cannot rotate **any** set:  $T_{ij} \rightarrow e^{i\theta_{ij}} T_{ij} \quad \theta_{ij} \in \mathbb{R}$
  - ★ consequence of linear 'LHS'  $\propto$  quadratic 'RHS'
3. Unitary parametrizations of data provide constraints that experiment may violate
  - ★ *normalization*, in particular
  - ★ next slide:  $^{17}\text{O}$  compound system

# R-matrix formalism

INTERIOR (Many-Body) REGION  
(Microscopic Calculations)



ASYMPTOTIC REGION  
(S-matrix, phase shifts, etc.)

$$(r_{c'})|\psi_c^+\rangle = -I_{c'}(r_{c'})\delta_{c'c} + O_{c'}(r_{c'})S_{c'c}$$

Measurements

$$\mathcal{L}_B = \sum_c |c\rangle(c| \left( \frac{\partial}{\partial r_c} r_c - B_c \right)),$$

$$(\mathbf{r}_c|c) = \frac{\hbar}{\sqrt{2\mu_c a_c}} \frac{\delta(r_c - a_c)}{r_c} [(\phi_{s_1}^{\mu_1} \otimes \phi_{s_2}^{\mu_2})_s^\mu \otimes Y_l^m(\hat{\mathbf{r}}_c)]_J^M$$

$$R_{c'c} = (c'| (H + \mathcal{L}_B - E)^{-1} | c) = \sum_\lambda \frac{(c'|\lambda)(\lambda|c)}{E_\lambda - E}$$

Bloch operator  $\mathcal{L}_B = \sum_c |c\rangle(c| \left[ \frac{\partial}{\partial r_c} r_c - B_c \right])$  ensures Hermiticity of Hamiltonian restricted to internal region

- R-matrix theory: **unitary**, multichannel parametrization of (not just resonance) data

- Interior/Exterior regions

- Interior: strong interactions
- Exterior: Coulomb/non-polarizing interactions
- Channel surface

$$\mathcal{S}_c : r_c = a_c \quad \mathcal{S} = \sum_c \mathcal{S}_c$$

- R-matrix elements

- Projections on channel surface functions  $(\mathbf{r}_c|c)$  of Green's function

$$G_B = [H + \mathcal{L}_B - E]^{-1}$$

- Boundary conditions

$$B_c = \frac{1}{u_c(a_c)} \frac{du_c}{dr_c} \Big|_{r_c=a_c}$$

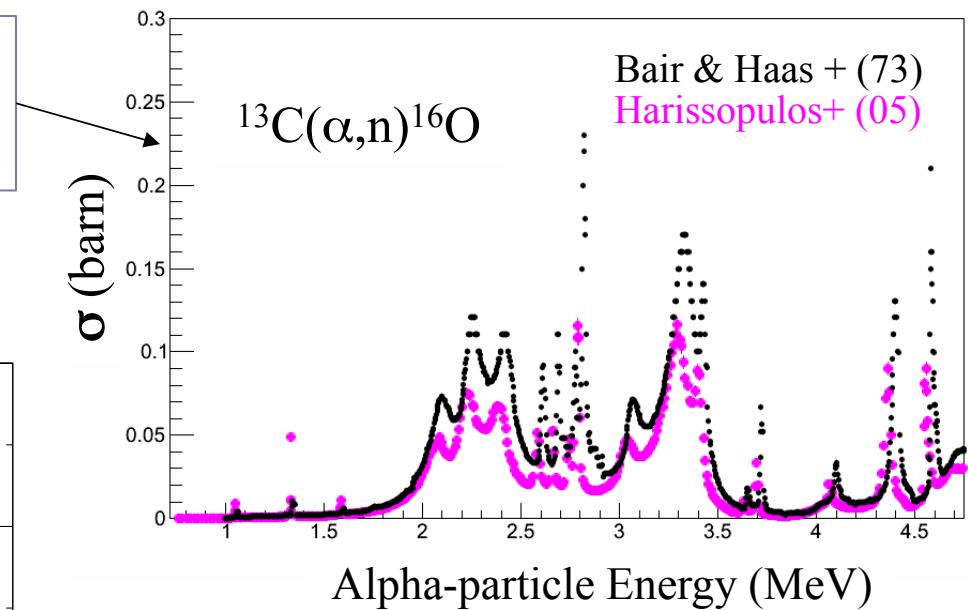
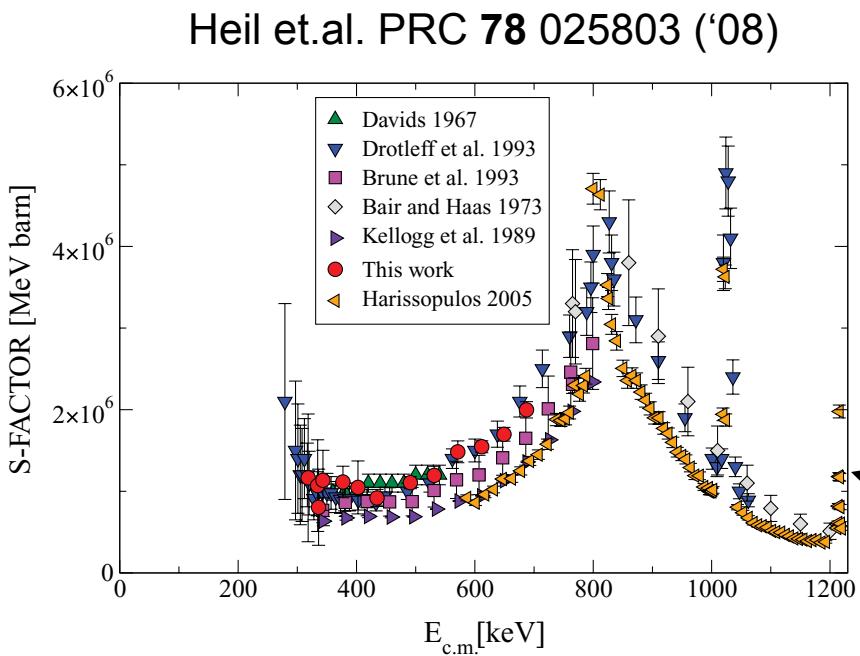
# <sup>17</sup>O analysis configuration

Channel	$a_c$ (fm)	$I_{\max}$
n+ <sup>16</sup> O	4.3	4
$\alpha$ + <sup>13</sup> C	5.4	5

Reaction	Energies (MeV)	# data points	Data types
<sup>16</sup> O(n,n) <sup>16</sup> O	$E_n = 0 - 7$	2718	$\sigma_T$ , $\sigma(\theta)$ , $P_n(\theta)$
<sup>16</sup> O(n, $\alpha$ ) <sup>13</sup> C	$E_n = 2.35 - 5$	850	$\sigma_{int}$ , $\sigma(\theta)$ , $A_n(\theta)$
<sup>13</sup> C( $\alpha$ ,n) <sup>16</sup> O	$E_\alpha = 0 - 5.4$	874	$\sigma_{int}$
<sup>13</sup> C( $\alpha$ , $\alpha$ ) <sup>13</sup> C	$E_\alpha = 2 - 5.7$	1296	$\sigma(\theta)$
total		5738	8

# $^{17}\text{O}$ compound system: experimental status

Recent (Harissopoulos '05)  
measurement  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  vs. older  
(Bair & Haas '73)



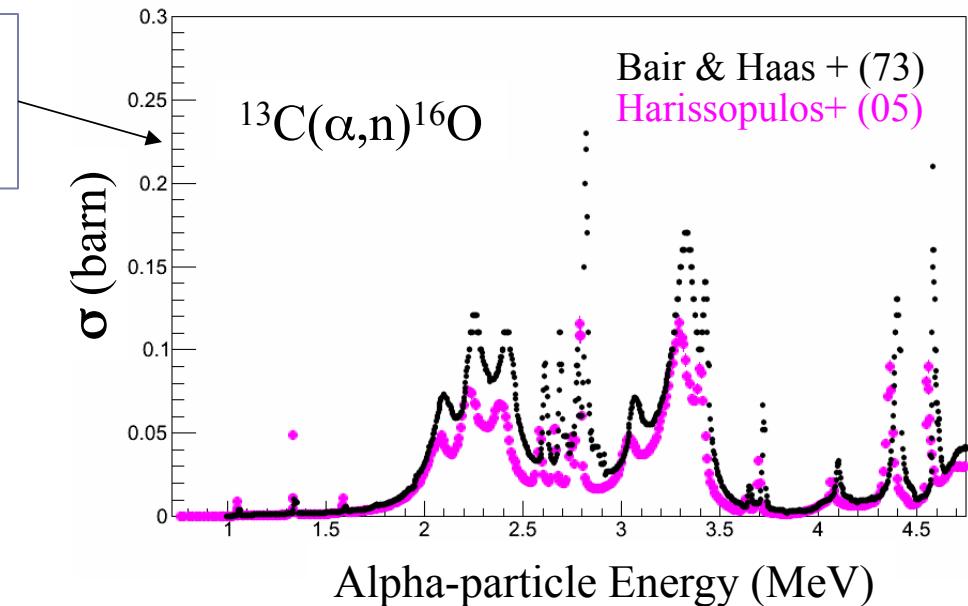
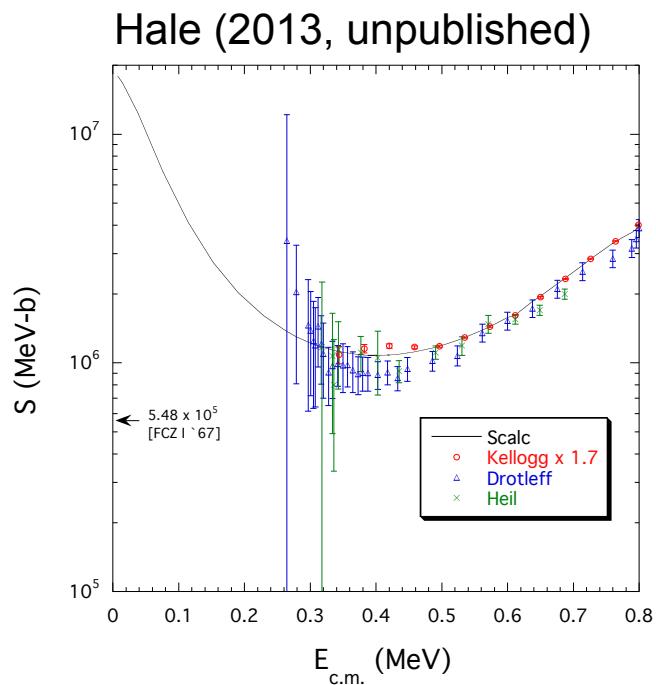
Credit: S. Kunieda

- Harissopoulos(05) data 2/3\*B&H(73)
- Heil(08) data consistent with B&H

Tempting to conclude that B&H73 was right all along!

# $^{17}\text{O}$ compound system: experimental status

Recent (Harissopoulos '05)  
measurement  $^{13}\text{C}(\alpha, n)^{16}\text{O}$  vs. older  
(Bair & Haas '73)



Credit: S. Kunieda

- Subthreshold  $\frac{1}{2}^+$
- deep min in  $\sigma_T$
- $S(0) \gg S_{\text{FCZ67}}(0)$

Tempting to conclude that B&H73 was right all along!

# R-matrix analyses support B&H73/Heil08

## ■ LANL R-matrix fit to Bair & Haas '73

→ two-channel fit:  $(^{16}\text{O},\text{n})$  &  $(^{13}\text{C},\alpha)$

$$\cdot \ell_n = 0, \dots, 4; \quad \ell_\alpha = 0, \dots, 5$$

→ data included:  $\sigma_T(E)$

•  $^{16}\text{O}(\text{n},\text{n})$ ,  $^{16}\text{O}(\text{n},\alpha)$ ,  $^{13}\text{C}(\alpha,\text{n})$

$$\cdot \sigma_{el}, d\sigma/d\Omega, A_y$$

•  $\chi^2$  minimization: normalizations float

→ Test Hariss05 data

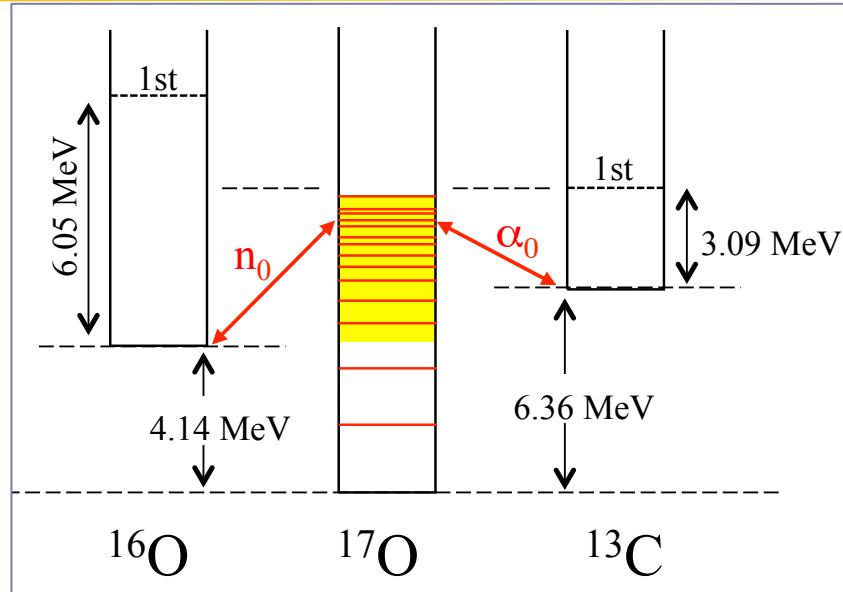
• remove B&H73/Heil08 data

• fix Hariss05 norm to unity

• unable to obtain fit with realistic  $\chi^2 (< 2.0)$

• now allow Hariss05 norm to float

• requires scale factor of  $\sim 1.5$ , consistent with B&H73



Credit: S. Kunieda

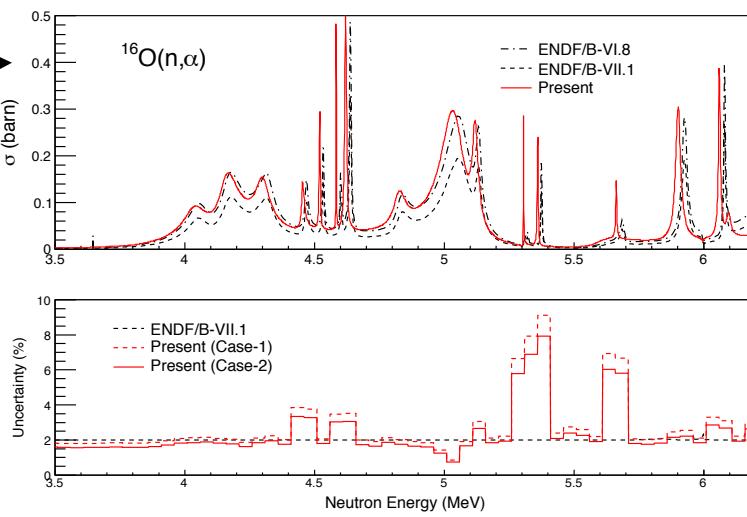
## ■ Kunieda/Kawano analysis [ND2013]

→ similar to LANL R-matrix(EDA)/ENDF/B-VI.8

→ with independent R-matrix code

→ KK give uncertainty analysis: see ND2013 proceedings in *Nucl. Data Sh.*

→ **Right to conclude B&H73 data correct on the basis of unitarity!**



Slide 15

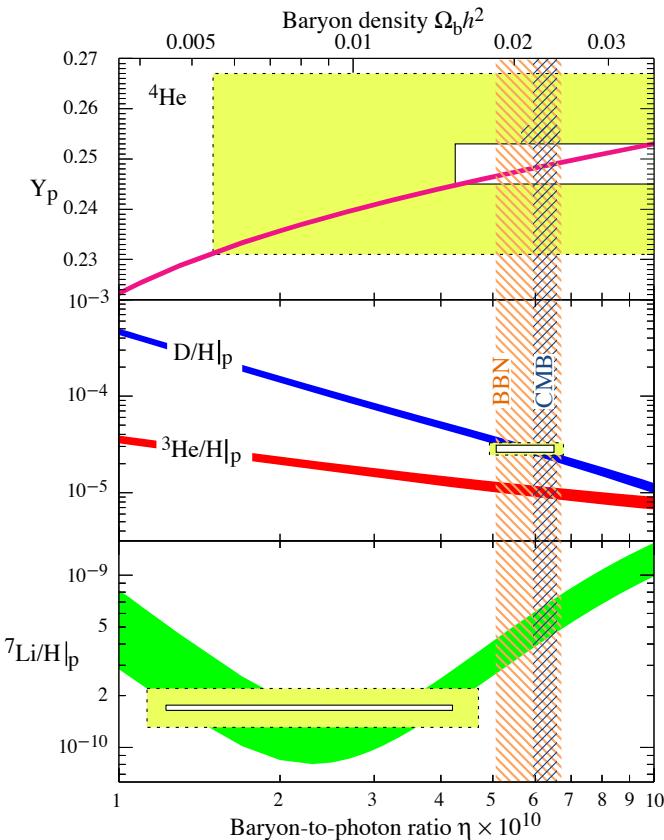
# Toward a unitary reaction network for BBN

## ■ Primordial nucleosynthesis

- Can unitarity play a role in precision BBN?
- D,  $^4\text{He}$  abundances agree with theo/expl uncertainties
- At  $\eta_{\text{wmap}}$  (CMB)  ${}^7\text{Li}/\text{H}|_{\text{BBN}} \sim (2.2\text{--}4.2) * {}^7\text{Li}/\text{H}|_{\text{halo}}$ \*
- Discrepancy  $\sim 4.5\text{--}5.5\sigma$  → the “Li problem”

## ■ Resonant destruction ${}^7\text{Li}$

- Prod. mass 7 “well understood”; destruction not
- Cyburt & Pospelov *arXiv:0906.4373; IJMPE, 21(2012)*
  - ${}^7\text{Be}(\text{d},\text{p})\alpha\alpha$  &  ${}^7\text{Be}(\text{d},\gamma){}^9\text{B}$  resonant enhancement
  - Identify  ${}^9\text{B}$   $E_{5/2^+} \approx 16.7$  MeV  $\approx E_{\text{thr}}(\text{d}+{}^7\text{Be}) + 200$  keV
    - *Near threshold*
  - $(E_r, \Gamma_d) \approx (170\text{--}220, 10\text{--}40)$  keV solve Li problem
- ‘Large’ widths
  - Conclude “large channel radius” required



**NB:** both approaches assume validity of TUNL-NDG tables

# <sup>9</sup>B analysis: included data

- **<sup>6</sup>Li+<sup>3</sup>He elastic** *Buzhinski et.al., Izv. Rossiiskoi Akademii Nauk, Ser.Fiz., Vol.43, p.158 (1979)*
  - Differential cross section
  - $1.30 \text{ MeV} < E(^3\text{He}) < 1.97 \text{ MeV}$
- **<sup>6</sup>Li+<sup>3</sup>He → p+<sup>8</sup>Be\*** *Elwyn et.al., Phys. Rev. C 22, 1406 (1980)*
  - Integrated cross section
  - Quasi-two-body, excited-state averaged final channel
  - $0.66 \text{ MeV} < E(^3\text{He}) < 5.00 \text{ MeV}$
- **<sup>6</sup>Li+<sup>3</sup>He → d+<sup>7</sup>Be** *D.W. Barr & J.S. Gilmore, unpublished (1965)*
  - Integrated cross section
  - $0.42 \text{ MeV} < E(^3\text{He}) < 4.94 \text{ MeV}$
- **<sup>6</sup>Li+<sup>3</sup>He → γ+<sup>9</sup>B** *Aleksic & Popic, Fizika 10, 273-278 (1978)*
  - Integrated cross section
  - $0.7 \text{ MeV} < E(^3\text{He}) < 0.825 \text{ MeV}$
  - New to <sup>9</sup>B analysis
- **New evaluation**
  - Separate <sup>8</sup>Be\* states
    - $2^+@200 \text{ keV}$  [16.9 MeV],  $1^+@650 \text{ keV}$  [17.6 MeV],  $1^+@1.1 \text{ MeV}$  [18.2 MeV]
  - n+<sup>8</sup>B:  $E_{\text{thresh}}(^3\text{He}) = 3 \text{ MeV}$
  - Simultaneous analysis with <sup>9</sup>Be mirror system

All data from  
**EXFOR/CSISRS**  
database (in  
C4 format)

# R-matrix configuration in EDA code

## Hadronic channels (in blue, not included)

$A_1 A_2 \pi$	$^3\text{He}^6\text{Li}^+(1)$		$p^8\text{Be}^{*+}(2)$		$d^7\text{Be}^-(3)$			
$\ell$	$S$	$\frac{3}{2}$	$\frac{1}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{1}{2}$
0		$^4S_{3/2}$	$^2S_{1/2}$	$^6S_{5/2}$	$^4S_{3/2}$	$^6S_{5/2}$	$^4S_{3/2}$	$^2S_{1/2}$
1		$^4P_{5/2,3/2,1/2}$	$^2P_{3/2,1/2}$	$^6P_{7/2,5/2,3/2}$	$^4P_{5/2,3/2,1/2}$	$^6P_{7/2,5/2,3/2}$	$^4P_{5/2,3/2,1/2}$	$^2P_{3/2,1/2}$
2		$^4D_{7/2,5/2,3/2,1/2}$	$^2D_{5/2,3/2}$	$^6D_{9/2,7/2,5/2,3/2,1/2}$	$^4D_{7/2,5/2,3/2,1/2}$	$^6D_{9/2,7/2,5/2,3/2,1/2}$	$^4D_{7/2,5/2,3/2,1/2}$	$^2D_{5/2,3/2}$
	$E_{\text{thr}}(\text{CM, MeV})$	16.6			16.7		16.5	

Electromagnetic channel:  $\gamma + {}^9B \rightarrow E_1^{3/2}, M_1^{5/2}, M_1^{3/2}, M_1^{1/2}, E_1^{5/2}, E_1^{1/2}$

Full model space:  
state number;  
channel pair;  
LS; J; channel  
radius [fm]

1	1	4s	3/2	7.50000000f	20	1	4p	1/2	7.50000000f
2	1	4d	3/2	7.50000000f	21	1	2p	1/2	7.50000000f
3	1	2d	3/2	7.50000000f	22	2	4p	1/2	5.50000000f
4	2	4s	3/2	5.50000000f	23	3	2s	1/2	7.00000000f
5	3	6p	3/2	7.00000000f	24	4	M1	1/2	50.00000000f
6	3	4p	3/2	7.00000000f	25	1	4d	7/2	7.50000000f
7	3	2p	3/2	7.00000000f	26	3	6p	7/2	7.00000000f
8	4	E1	3/2	50.00000000f	27	1	4d	5/2	7.50000000f
9	1	4p	5/2	7.50000000f	28	1	2d	5/2	7.50000000f
10	2	6p	5/2	5.50000000f	29	2	6s	5/2	5.50000000f
11	2	4p	5/2	5.50000000f	30	3	6p	5/2	7.00000000f
12	3	6s	5/2	7.00000000f	31	3	4p	5/2	7.00000000f
13	4	M1	5/2	50.00000000f	32	4	E1	5/2	50.00000000f
14	1	4p	3/2	7.50000000f	33	1	4d	1/2	7.50000000f
15	1	2p	3/2	7.50000000f	34	1	2s	1/2	7.50000000f
16	2	6p	3/2	5.50000000f	35	3	4p	1/2	7.00000000f
17	2	4p	3/2	5.50000000f	36	3	2p	1/2	7.00000000f
18	3	4s	3/2	7.00000000f	37	4	E1	1/2	50.00000000f
19	4	M1	3/2	50.00000000f	38	2	6p	7/2	5.50000000f

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Slide 18

# Analysis result: resonance structure

Ex(MeV)	Jpi	Gamma(keV)	Er(MeV)	ImEr(MeV)	E( <sup>3</sup> He)	Strength
16.46539	1/2-	768.46	- .1369	-0.3842	-0.2054	0.06 weak
17.11317	1/2-	0.14	0.5109	-0.6771E-04	0.7664	1.00 strong
17.20115	5/2-	871.63	0.5989	-0.4358	0.8984	0.40 weak
17.28086	3/2-	147.78	0.6785	-0.0739	1.0178	0.77 strong
17.66538	5/2+	33.33	1.0631	-0.0167	1.5947	0.98 strong
17.83619	7/2+	2036.21	1.2339	-1.0181	1.8509	0.15 weak
17.84773	3/2-	42.52	1.2454	-0.0213	1.8681	0.97 strong
18.04821	3/2+	767.11	1.4459	-0.3836	2.1689	0.54 weak
18.42292	1/2+	5446.32	1.8206	-2.7232	2.7309	0.03 weak
18.67716	1/2-	10278.41	2.0749	-5.1392	3.1124	0.15 weak
19.60923	3/2-	1478.22	3.0069	-0.7391	4.5104	0.52 weak

TUNL-NDG/ENSDF  
parameters

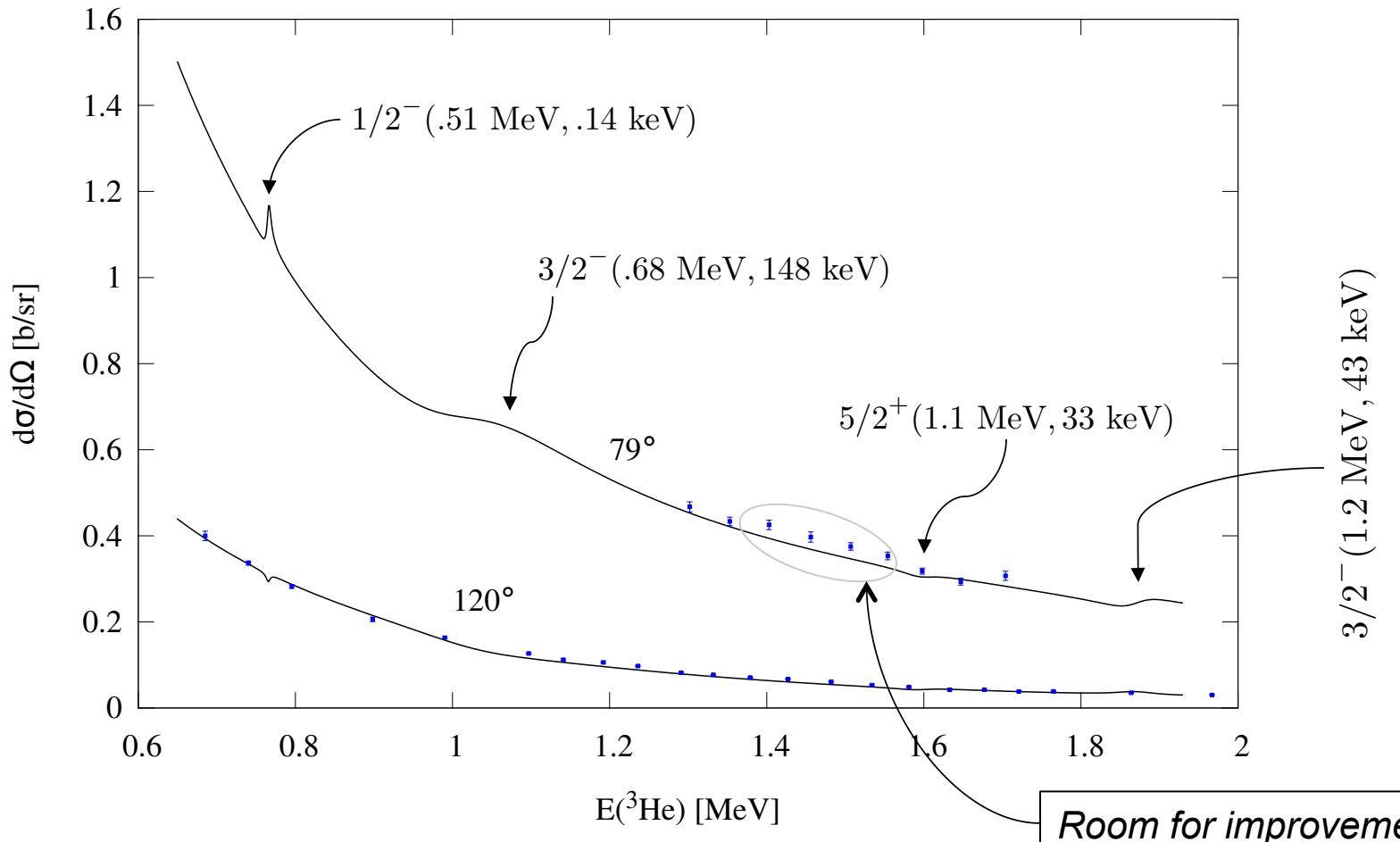
**NB: no strong resonance seen  
~100 keV of <sup>3</sup>He+<sup>6</sup>Li threshold**

$E_x^a$ (MeV ± keV)	$J^\pi; T$	$\Gamma_{c.m.}$ (keV)	Decay
$16.024 \pm 25$	$T = (\frac{1}{2})$	$180 \pm 16$	
$16.71 \pm 100^h$	$(\frac{5}{2}^+); (\frac{1}{2})$		
$17.076 \pm 4$	$\frac{1}{2}^-; \frac{3}{2}$	$22 \pm 5$	$(\gamma, {}^3\text{He})$
$17.190 \pm 25$		$120 \pm 40$	$p, d, {}^3\text{He}$
$17.54 \pm 100^{h,i}$	$(\frac{7}{2}^+); (\frac{1}{2})$		
$17.637 \pm 10^i$		$71 \pm 8$	$p, d, {}^3\text{He}, \alpha$

# Observable fit: ${}^3\text{He} + {}^6\text{Li}$ elastic DCS

${}^6\text{Li}({}^3\text{He}, \text{Elastic})$   
Differential cross section

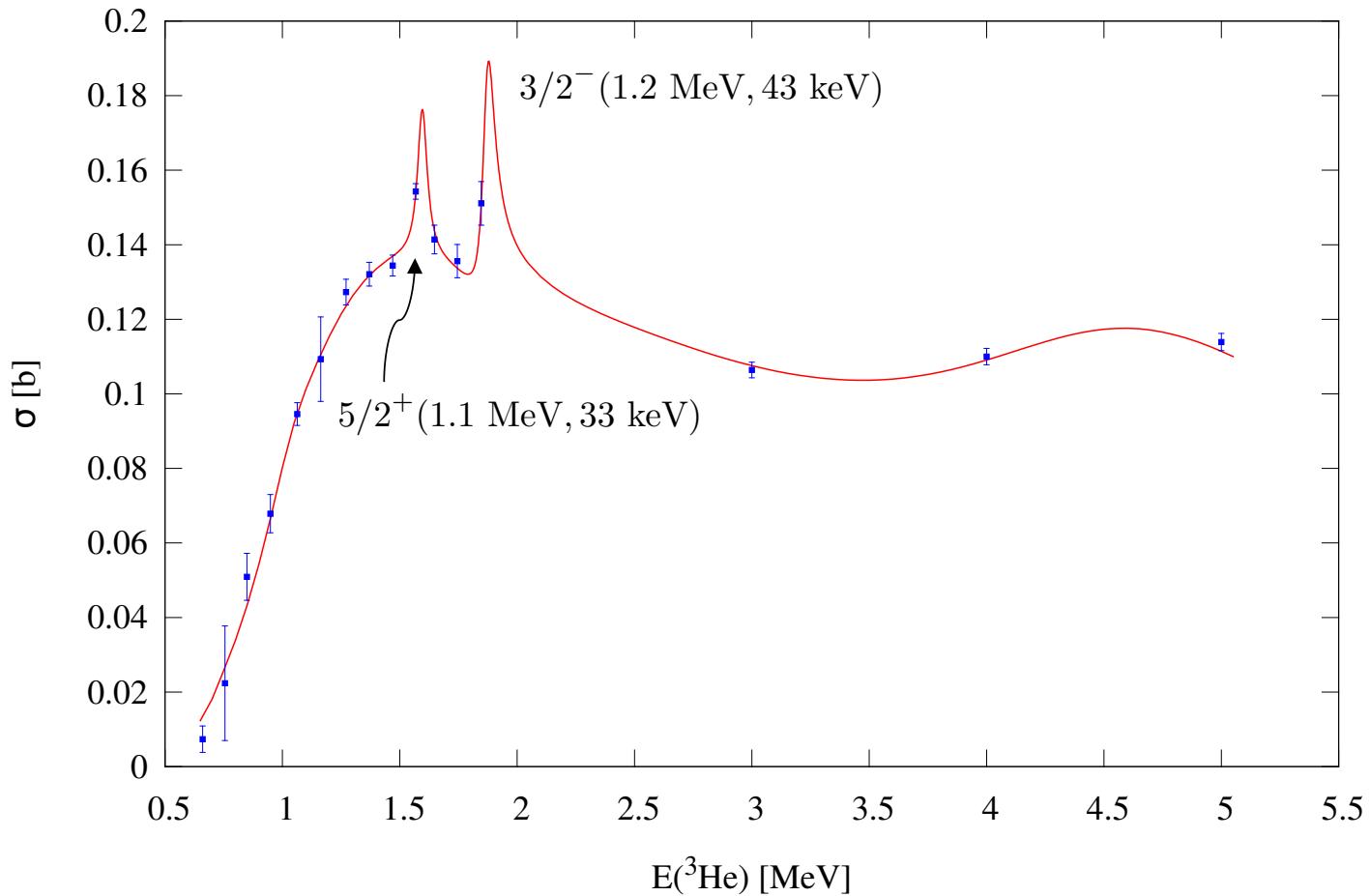
$\chi^2/N_{\text{data}} = 1.91$



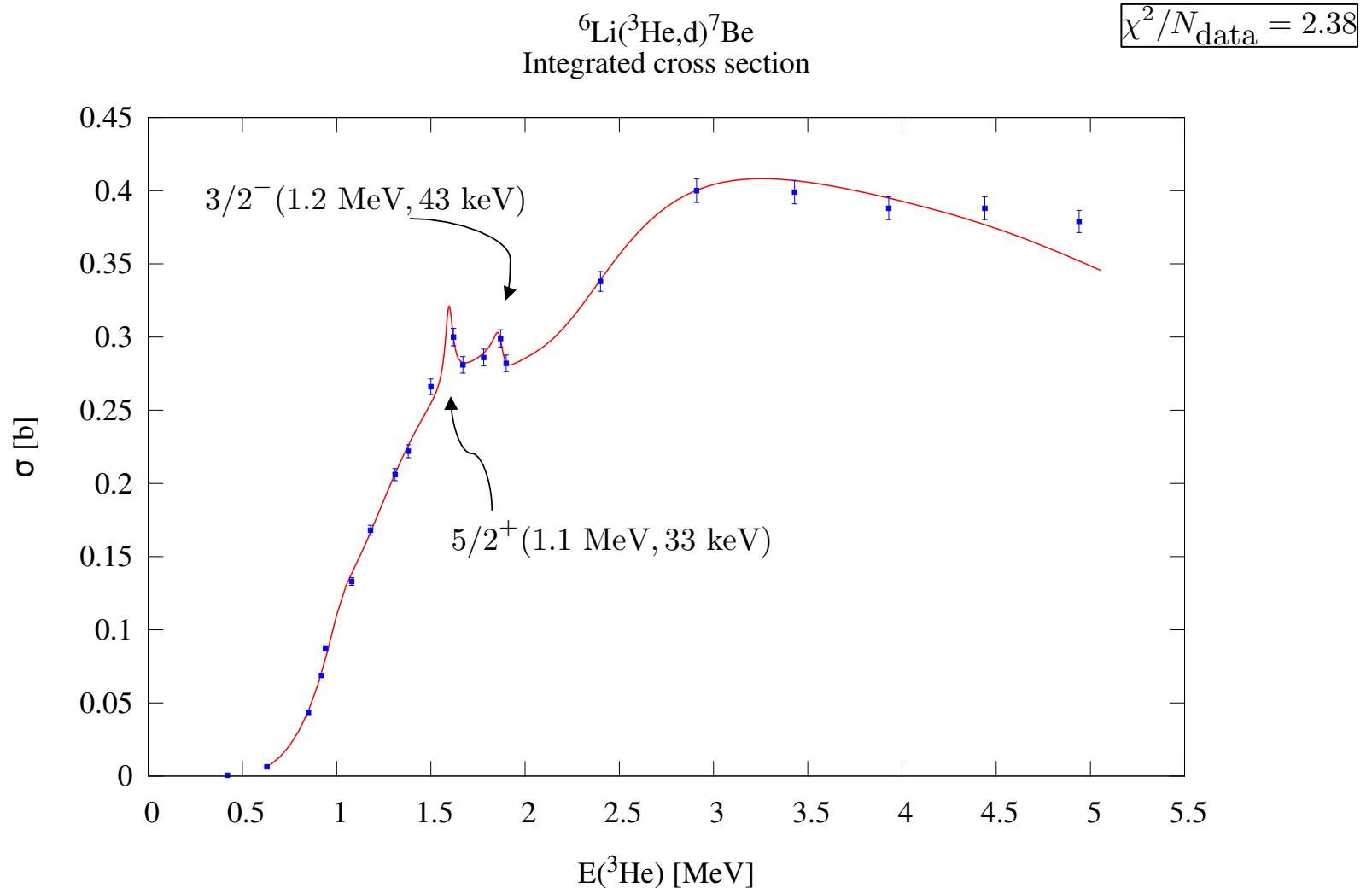
# Observable fit: ${}^6\text{Li}({}^3\text{He},\text{p}){}^8\text{Be}^*$ integrated x-sec

${}^6\text{Li}({}^3\text{He},\text{p}){}^8\text{Be}^*$   
Integrated cross section

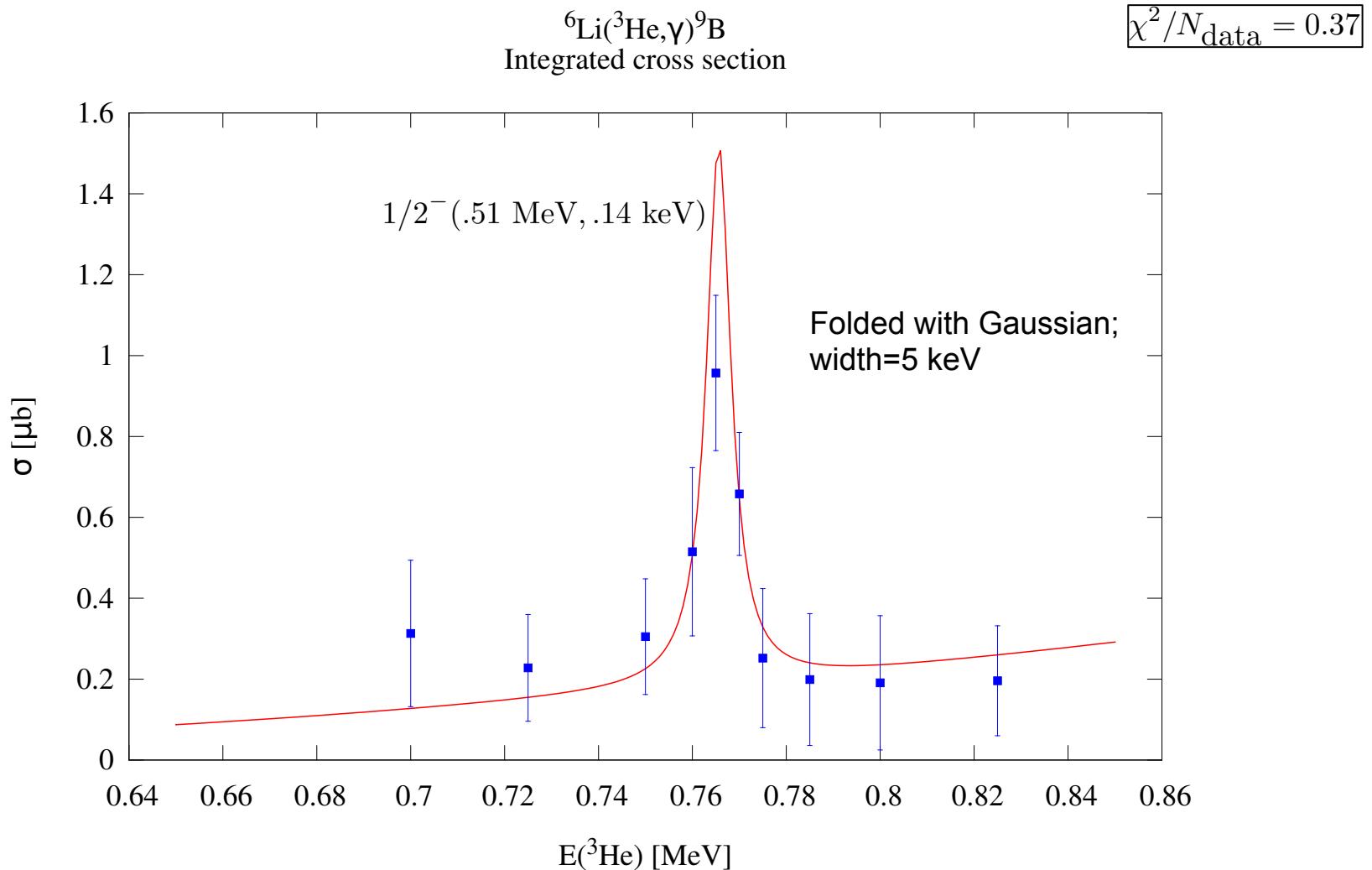
$\chi^2/N_{\text{data}} = 0.55$



# Observable fit: ${}^6\text{Li}({}^3\text{He},\text{d}){}^7\text{Be}$ integrated x-sec



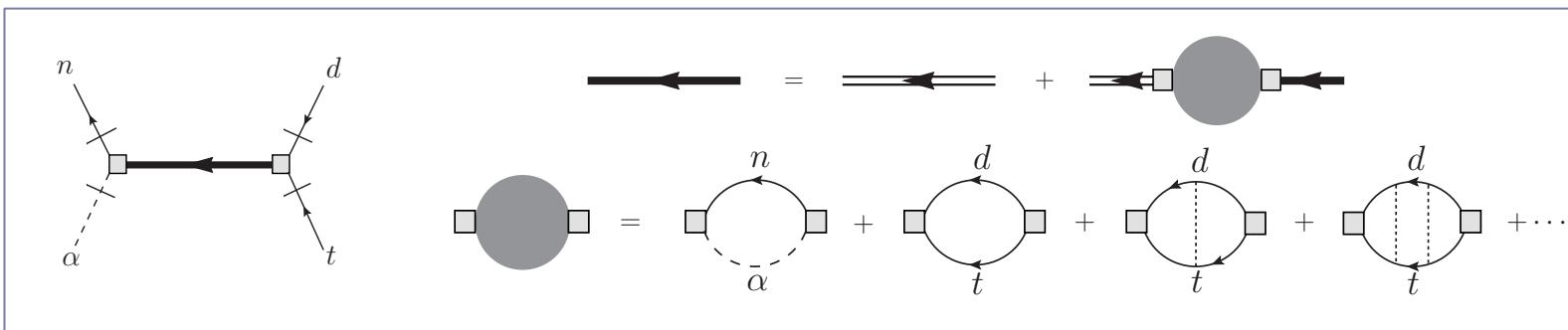
# Observable fit: ${}^6\text{Li}({}^3\text{He},\gamma){}^9\text{B}$ integrated x-sec



# Effective field theory $\leftrightarrow$ R matrix: dt $\rightarrow$ n $\alpha$

- Exactly soluble EFT with ‘wrong-sign’ free Lagrangian and DOF:

Particle	Spin	Operators	Mass	Binding
Alpha	$0^+$	$\phi_\alpha^\dagger(\mathbf{r}, t), \phi_\alpha(\mathbf{r}, t)$	$m_\alpha = 2m_p + 2m_n$	$\epsilon_\alpha$
Deuteron	$1^+$	$\phi_d^\dagger(\mathbf{r}, t), \phi_d(\mathbf{r}, t)$	$m_d = m_p + m_n$	$\epsilon_d$
Neutron	$\frac{1}{2}^+$	$\psi_n^\dagger(\mathbf{r}, t), \psi_n(\mathbf{r}, t)$	$m_n$	$\epsilon_n \equiv 0$
Triton	$\frac{1}{2}^+$	$\psi_t^\dagger(\mathbf{r}, t), \psi_t(\mathbf{r}, t)$	$m_t = m_p + 2m_n$	$\epsilon_t$
${}^5\text{He}^*$	$\frac{3}{2}^+$	$\psi_*^\dagger(\mathbf{r}, t), \psi_*(\mathbf{r}, t)$	$m_* = 2m_p + 3m_n$	$\epsilon_*$



$$\sigma_{dt \rightarrow n\alpha} = \frac{8}{9} 4\pi m_{n\alpha} \frac{p_{n\alpha}^5}{v_{dt}} \frac{g_{dt}^2}{4\pi} \frac{g_{n\alpha}^2}{4\pi} \left| \psi_{\mathbf{P}_{dt}}^{(0)}(0) \right|^2 \left[ \left( \frac{p_{dt}^2}{2m_{dt}} - E_* - \frac{g_{dt}^2}{4\pi} \Delta(W) \right)^2 + \left[ \frac{g_{dt}^2}{4\pi} 2m_{dt} p_{dt} \left| \psi_{\mathbf{P}_{dt}}^{(C)}(0) \right|^2 + \frac{g_{n\alpha}^2}{4\pi} \frac{2}{3} m_{n\alpha} p_{n\alpha}^5 \right]^2 \right]$$

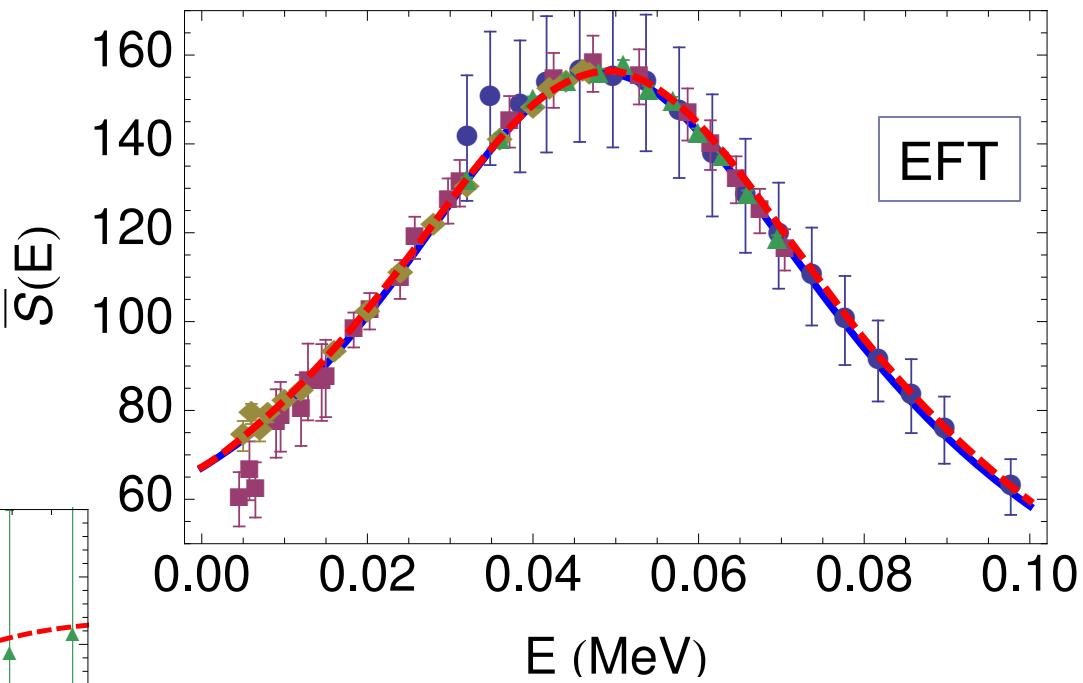
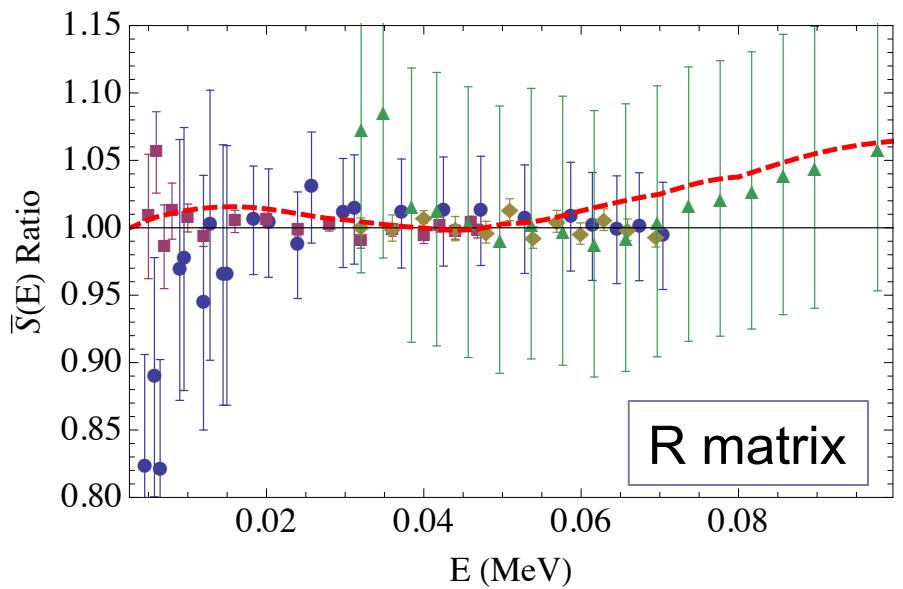
- Now consider R-matrix (d,t)&(n, $\alpha$ ) in the limit  $a_d, a_n \rightarrow 0$

$$\gamma_d^2 = -\frac{g_d^2}{2\pi} \frac{\mu_d}{\hbar^2 a_d} \quad \text{and} \quad \gamma_n^2 = -\frac{g_n^2}{6\pi} \frac{\mu_n}{\hbar^2 a_n^5}$$

$$\sigma_{n,d}^{3/2+} = \frac{32}{9\hbar v_d} \frac{g_d^2}{4\pi} \frac{g_n^2}{4\pi} \frac{\mu_n}{\hbar^2} k_n^5 C_0^2(\eta_d) \left| E - E_\lambda - \Delta_d(E) - i \left[ \frac{g_d^2}{2\pi} \frac{\mu_d}{\hbar^2} k_d C_0^2(\eta_d) + \frac{g_n^2}{6\pi} \frac{\mu_n}{\hbar^2} k_n^5 \right] \right|^{-2}$$

# Effective field theory $\leftrightarrow$ R matrix: $dt \rightarrow n\alpha$

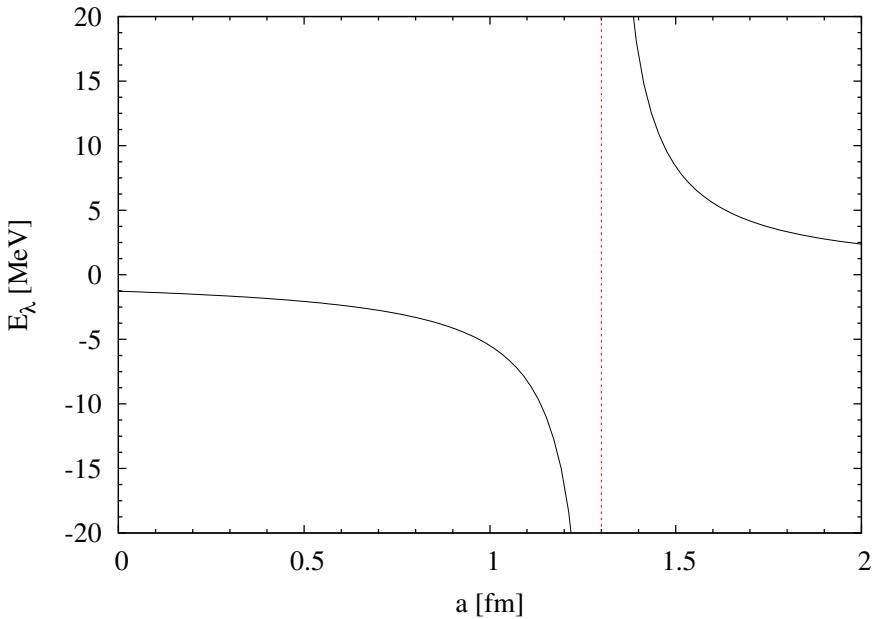
- Dim'less astrophys factor vs  $E_d(\text{CM})$
- Solid-blue: 3 par EFT fit  $\chi^2/\text{dof} \sim 0.8$
- Dashed-red: Bosch & Hale (1992)
- 2665 data/117 pars/ $\chi^2/\text{dof} \sim 1.6$



- Three parameter fit
- Unphysical reduced widths:  $\gamma_c^2 < 0$

# Zero channel radius limit: $^1S_0$ $np$ scattering

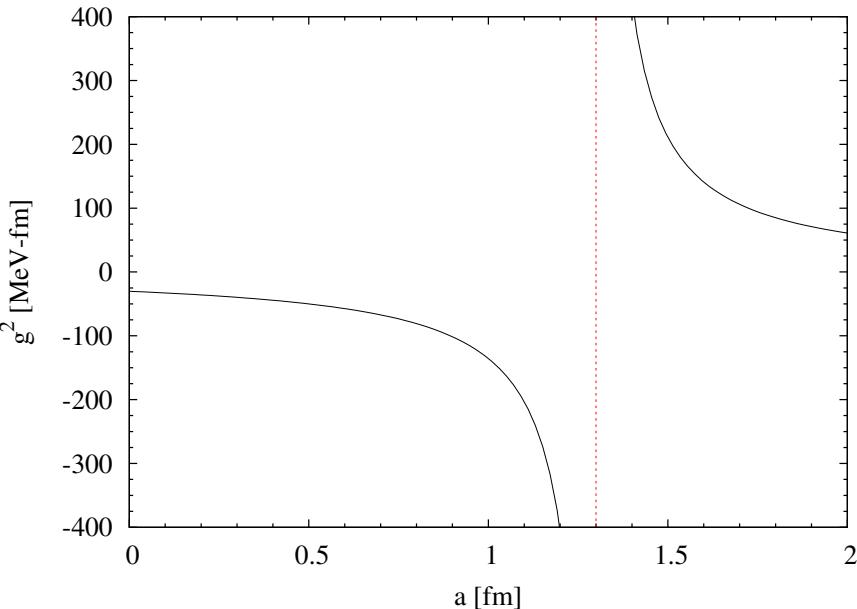
$$k \cot \delta_0(E) = \frac{E_\lambda - \frac{\hbar^2 k^2}{2\mu} + kg^2 \tan ka}{g^2 - (E_\lambda - \frac{\hbar^2 k^2}{2\mu}) \frac{1}{k} \tan ka}$$



$$E_\lambda(a) = \frac{\hbar^2(a_0 - a)}{2\mu[r_0a_0^2/2 - a^3/3 - aa_0(a_0 - a)]}$$

**Pole position:**  $a_p = a_0 + \left\{ a_0^3 \left[ \frac{3r_0}{2a_0} - 1 \right] \right\}^{1/3}$

$$\begin{aligned} g^2(a) &= (a - a_0)E_\lambda(a) \\ &= -\frac{\hbar^2(a_0 - a)^2}{2\mu[r_0a_0^2/2 - a^3/3 - aa_0(a_0 - a)]} \end{aligned}$$



# Summary, findings & future work

## Summary/findings

- Provided overview of current work in the LANL light nuclear reaction program
- Emphasize the utility of multichannel, unitary parametrization of light nuc data
  - $^{17}\text{O}$  norm issue
  - $^9\text{B}$  resonance spectrum: no resonances in  $^9\text{B}$  that reside within  $\sim 200$  ( $\sim 100$ ) keV of the  $d+^7\text{Be}$  ( $^3\text{He}+^6\text{Li}$ ) threshold with ‘large’ widths 10—40 keV
  - Appears to rule out scenarios considered by *Cyburt & Pospelov (2009)* that low-lying, robust resonance in  $^9\text{B}$  could explain the “Li problem”

## Near-term, Future Work

- Complete  $^{13,14}\text{C}$  analyses
- $NN$  up to 200 MeV
- Improvements in the  $^9\text{B}$  analysis: more channels; incorporate  $p+^8\text{Be}^*$  angular data; proper treatment three-body final states
- Extend EFT—R-matrix approach to multichannel, multilevel problems

# Supplementary material

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Additional slides follow

# Motivation

## ■ Cross section evaluation & resonance structure

→ *Nucl. Phys. A745, 155, 2004(2011)*

$E_x^a$ (MeV ± keV)	$J^\pi; T$	$\Gamma_{c.m.}$ (keV)	Decay
$16.024 \pm 25$	$T = (\frac{1}{2})$	$180 \pm 16$	
$16.71 \pm 100^h$	$(\frac{5}{2}^+); (\frac{1}{2})$		
$17.076 \pm 4$	$\frac{1}{2}^-; \frac{3}{2}$	$22 \pm 5$	$(\gamma, {}^3\text{He})$
$17.190 \pm 25$		$120 \pm 40$	$p, d, {}^3\text{He}$
$17.54 \pm 100^{h,i}$	$(\frac{7}{2}^+); (\frac{1}{2})$		
$17.637 \pm 10^i$		$71 \pm 8$	$p, d, {}^3\text{He}, \alpha$

## ■ Astrophysical applications

→ Big bang nucleosynthesis

- Nuclear physics solution to  ${}^7\text{Li}$  predicted overproduction problem? (cf. Hoyle)
- Details next slide.

## ■ Purpose within Los Alamos Nat. Lab programmatic

→ Continue the R-matrix program for various end-users

→ Ongoing/upcoming analysis releases:  ${}^7\text{Be}$ ,  ${}^{13}\text{C}$  [G. Hale Tues. Session GA 2],  ${}^{14}\text{C}$ ,  ${}^{17}\text{O}$ , ...

# Implementation in EDA

- **EDA = Energy Dependent Analysis**

- Adjust  $E_\lambda$  &  $\gamma_{c\lambda}$

- **Any number of two-body channels**

- Arbitrary spins, masses, charges (incl. mass zero)

- **Scattering observables**

- Wolfenstein trace formalism

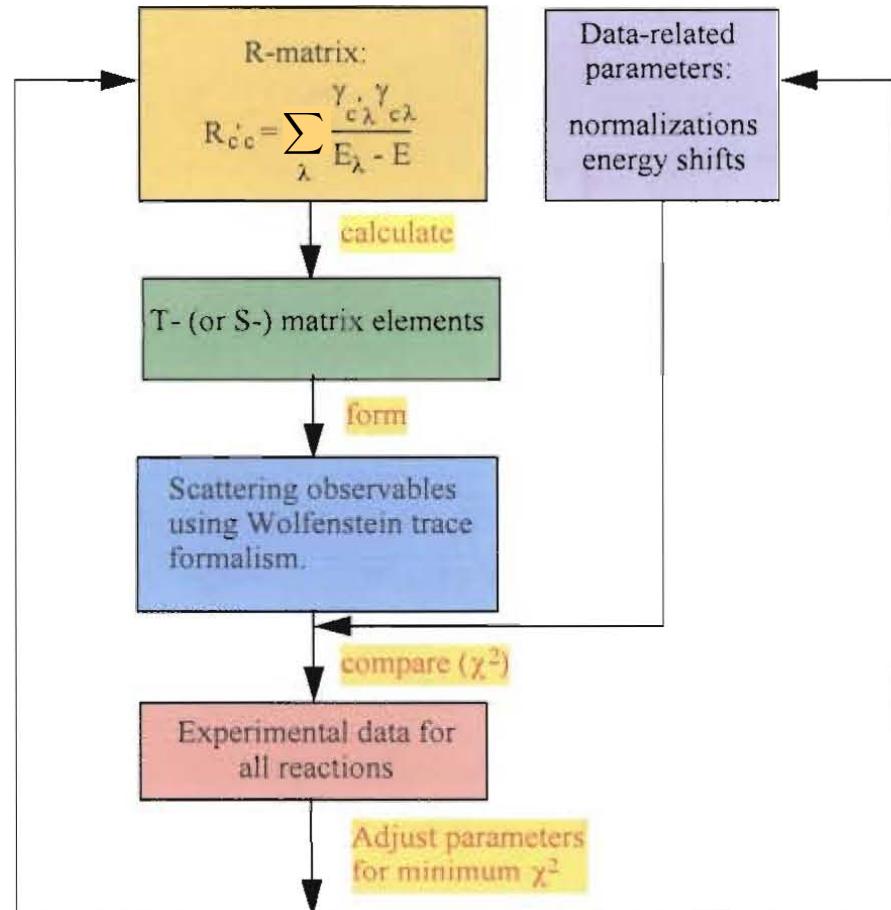
- **Data**

- Normalization
  - Energy shifts
  - Energy resolution/spread

- **Fit solution**

$$\chi^2_{EDA} = \sum_i \left[ \frac{nX_i(\mathbf{p}) - R_i}{\delta R_i} \right]^2 + \left[ \frac{nS - 1}{\delta S/S} \right]^2$$

- **Covariance determined**



# Electromagnetic channels

## ■ One-photon sector of Fock space

→ Photon ‘wave function’

$$\mathbf{A}_k(\mathbf{r}) = \left( \frac{2}{\pi \hbar c} \right)^{1/2} \sum_{jm} i^j \sum_{\lambda', \lambda=e,m,0} \mathbf{Y}_{jm}^{(\lambda')}(\hat{\mathbf{r}}) u_{\lambda' \lambda}^j(r) \mathbf{Y}_{jm}^{(\lambda)}(\hat{\mathbf{k}}) \cdot \chi$$

→ Radial part

$$u_{ee}^j = -[f'_j(\rho) + t_{ee}^j h_j^+(\rho)] \quad u_{0e}^j = -\frac{\sqrt{j(j+1)}}{\rho} [f_j(\rho) + t_{e0}^j h_j^+(\rho)]$$
$$u_{mm}^j = [f_j(\rho) + t_{mm}^j h_j^+(\rho)] \quad u_{0m}^j = u_{me}^j = u_{em}^j = 0$$

→ Photon channel surface functions

$$(\mathbf{r}_c|c) = \left( \frac{\hbar c}{2\rho_\gamma} \right)^{1/2} \frac{\delta(r_\gamma - a_\gamma)}{r_\gamma} \left[ \phi_{s\nu} \otimes \mathbf{Y}_{jm}^{(e,m)}(\hat{\mathbf{r}}_\gamma) \right]_{JM}$$

- Photon ‘mass’:  $\hbar k_\gamma/c$

→ R-matrix definition preserved

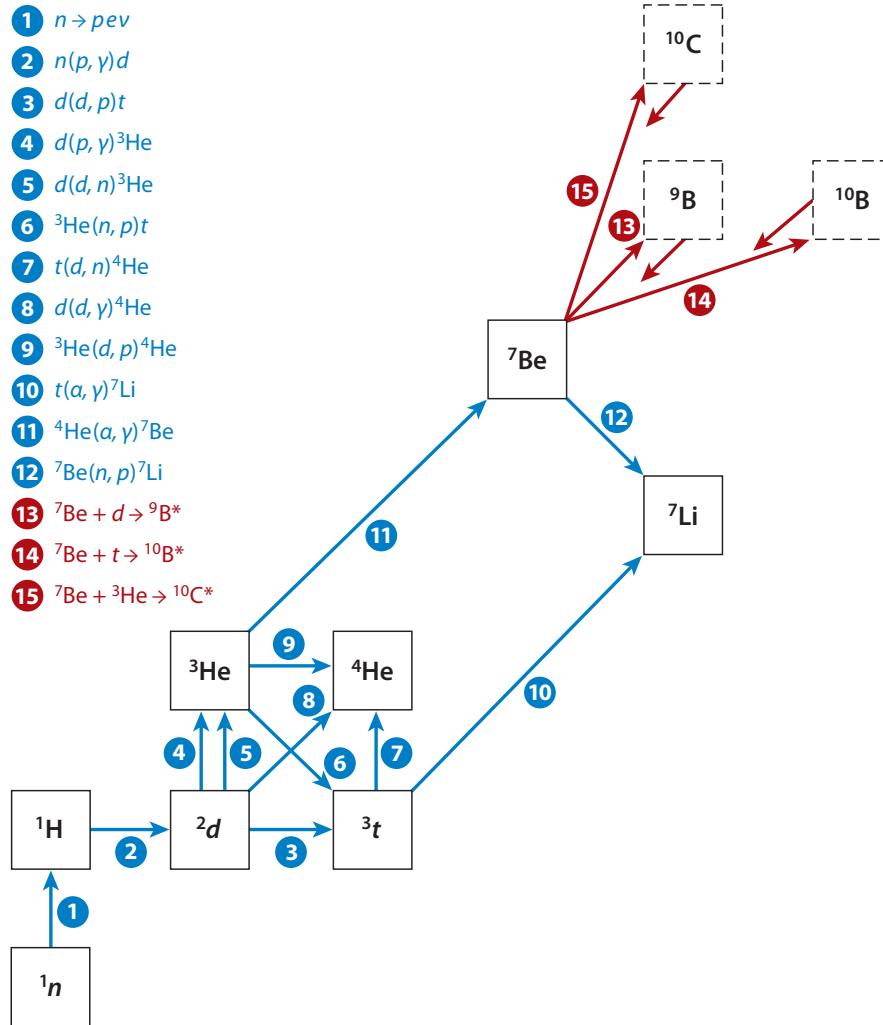
$$(c'|\psi) = \sum_c R_{c'c}^B (c| \frac{\partial}{\partial r_c} r_c - B_c | \psi)$$

$$\begin{aligned} \mathbf{T} &= \rho^{1/2} \mathbf{O}^{-1} \mathbf{R}_L \mathbf{O}^{-1} \rho^{1/2} - \mathbf{F} \mathbf{O}^{-1} \\ \mathbf{R}_L &= [\mathbf{R}_B^{-1} - \mathbf{L} + \mathbf{B}]^{-1} \\ \mathbf{L} &= \rho \mathbf{O}' \mathbf{O}^{-1} \\ \mathbf{F} &= \text{Im } \mathbf{O} \end{aligned}$$

# BBN reaction network (simplified)

■ Fields *Annu. Rev. Nucl. Part. Sci.* 2011. 61:47–68

- 1  $n \rightarrow pe\nu$
- 2  $n(p, \gamma)d$
- 3  $d(d, p)t$
- 4  $d(p, \gamma)^3\text{He}$
- 5  $d(d, n)^3\text{He}$
- 6  $^3\text{He}(n, p)t$
- 7  $t(d, n)^4\text{He}$
- 8  $d(d, \gamma)^4\text{He}$
- 9  $^3\text{He}(d, p)^4\text{He}$
- 10  $t(a, \gamma)^7\text{Li}$
- 11  $^4\text{He}(a, \gamma)^7\text{Be}$
- 12  $^7\text{Be}(n, p)^7\text{Li}$
- 13  $^7\text{Be} + d \rightarrow ^9\text{B}^*$
- 14  $^7\text{Be} + t \rightarrow ^{10}\text{B}^*$
- 15  $^7\text{Be} + ^3\text{He} \rightarrow ^{10}\text{C}^*$

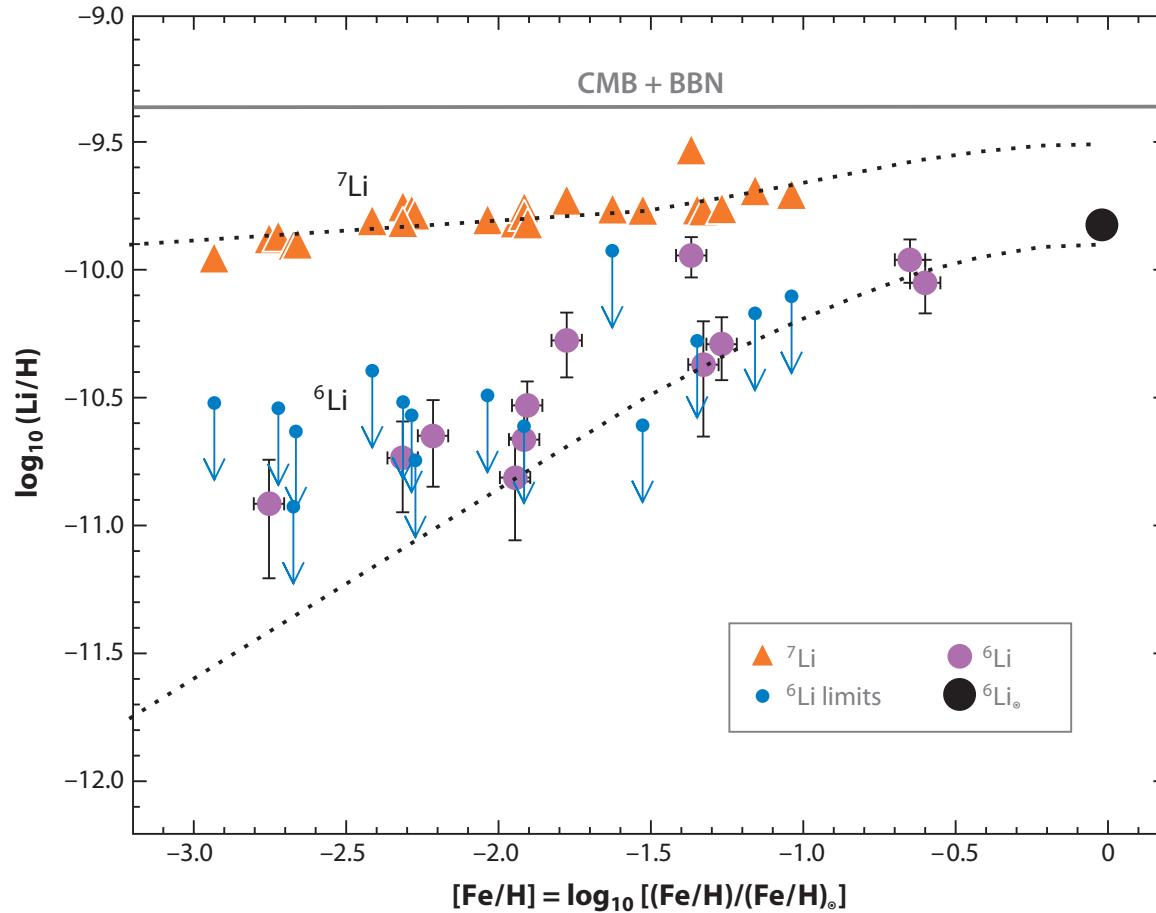


UNCLASSIFIED

Slide 32

# Spite Plateau

## ■ Measurement of primordial $^7\text{Li}$ from low-metallicity halo dwarf stars



Asplund M, et al. *Astrophys. J.* 644:229 (2006)

# Analysis result: resonance structure

Ex(MeV)	Jpi	Gamma(keV)	Er(MeV)	ImEr(MeV)	E( <sup>3</sup> He)	Strength
16.46539	1/2-	768.46	- .1369	-0.3842	-0.2054	0.06 weak
17.11317	1/2-	0.14	0.5109	-0.6771E-04	0.7664	1.00 strong
17.20115	5/2-	871.63	0.5989	-0.4358	0.8984	0.40 weak
17.28086	3/2-	147.78	0.6785	-0.0739	1.0178	0.77 strong
17.66538	5/2+	33.33	1.0631	-0.0167	1.5947	0.98 strong
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19.60923	3/2-	1478.22	3.0069	-0.7391	4.5104	0.52 weak

S-matrix pole & residue *Hale, Brown, Jarmie PRL 59 '87*

$$\mathcal{E}_{\lambda'\lambda} = E_\lambda \delta_{\lambda'\lambda} - \sum_c \gamma_{c\lambda'} [L_c(E) - B_c] \gamma_{c\lambda}$$

$$E_0 = E_r - i\Gamma/2 \quad \text{residue: } i\rho_0 \rho_0^T$$

**NB: no strong resonance seen  
~100 keV of <sup>3</sup>He+<sup>6</sup>Li threshold**

$$\text{strength} = \frac{1}{\Gamma} \rho_0^\dagger \rho_0 = \frac{1}{\Gamma} \sum_c \Gamma_c$$

$$\rho_{0c} = \left( \frac{2k_{0c}a_c}{N} \right)^{1/2} \mathcal{O}_c^{-1}(k_{0c}a_c) \sum_{\lambda} (\lambda|\mu_0)$$

$$N = \sum_{\lambda'\lambda} (\lambda|\mu_0)(\lambda'|\mu_0) \left[ \delta_{\lambda'\lambda} + \sum_c \gamma_{c\lambda'} \frac{\partial L_c}{\partial E} \Big|_{E=E_0} \gamma_{c\lambda} \right]$$

$$L_c = r_c \frac{\partial \mathcal{O}_c}{\partial r_c} \mathcal{O}_c^{-1} \Big|_{r_c=a_c}$$